LFA-tuned matrix-free multigrid for the elastic Helmholtz equations

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The acoustic Helmholtz equation in heterogeneous media

$$Hp = \rho \nabla \cdot \rho^{-1} \nabla p + \omega^2 \kappa^2 (1 - \gamma \iota) p = q$$

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The elastic Helmholtz equation in isotropic heterogeneous media

$$\mathcal{H}\vec{u} =
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abla \cdot \mu
abla)\vec{u} + \omega^2
ho(1 - \gamma \iota)\vec{u} = \vec{q}$$

Pressure and shear wave velocities: $V_{\rho} = \sqrt{\frac{\lambda + 2\mu}{\rho}}$ and $V_{s} = \sqrt{\frac{\mu}{\rho}}$.

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Boundary Conditions

• Neumann / Dirichlet: cause reflections.

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Boundary Conditions

- Neumann / Dirichlet: cause reflections.
- Absorbing alternatives: Sommerfeld, ABC, PML.

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Numerical dispersion





FD Discretizations for acoustic Helmholtz

- 4-th order [Harari & Turkel 1994], [Singer & Turkel 1998]
- 6-th order [Turkel et. al. 2013]
- Non-compact [Dastur & Liau 2020]
- Lower dispersion compact [Liu 2015], [Chen et. al. 2012]
- Rotated grids [Alghamiry et. al. 2022]
- Wavelength adaptive [Xu & Gao 2018]

Solvers for acoustic Helmholtz

- Domain decomposition [Gander & Zhang 2013], [Hu & Li 2016], [Tausa et. al. 2020], [Daia et. al. 2022], [Claeys, MS2], [Dolean, MS2], [Gong, MS2], [Rees, IP6]
- Shifted Laplacian multigrid [Erlangga et. al. 2006], [Umetani et. al 2009], [Chen et. al. 2012], [Dwarka, MS2], [Chen, MS1]
- FFT [Beylkin 2009], [Osnabrugge 2016], [Wang et. al. 2020]

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Finite difference discretization and available solvers

FD Discretizations for elastic Helmholtz

- 2nd order staggered (rotated grids) [Virieux 1986]
- Compact nodal [Sketl & Pratt 1998], [Gosselin-Cliche & Giroux 2014]
- 4th order staggered [Levander 1988]
- 4th order staggered with mass spreading [Li et. al. 2016]

[Gap]

Multigrid adaptations for elastic Helmholtz

- Shifted Laplacian "as is" [Airaksinen et. al. 2009]
- Shifted Laplacian using line-smoothers. [Rizzuti & Mulder 2016]

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Multigrid: iterative methods for elliptic PDEs [Brandt 1977]

- Smoothers (i.e. w-Jacobi, GS) reduce oscillatory error modes
- Coarse grid correction (CGC) reduce the smooth modes



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Left image from: "Why Multigrid Methods Are So Efficient" [Yavneh, 2006]

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Being matrix-free is not (only) about implementation, but about improving complexity [Pazner, IP2].

Local Fourier analysis

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LFA for multiplicative overlapping smoothers

- The amplification factor $\approx \frac{|\text{error after}|}{|\text{error before}|}$
- We need "middle". "before" and "after" are not enough!

[Sivaloganathan 1991], [Maclachlan & Oosterlee 2011], [Rodrigo et. al. 2016], [Treister & Y. 2024]



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Shifted Laplacian [Erlangga et. al. 2006]

Use an attenuated matrix $H_s = H + i\alpha\omega^2 M$ as a preconditioner.

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Our adaptations: a new variable $p = -(\lambda + \mu)\nabla \cdot \vec{u}$.



• MAC staggered grid finite differences discretization: [McKee et. al. 2008]



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• Cell-wise Vanka smoother: invert a 5×5 matrix at every cell. [Vanka 1986]



Multigrid for elastic Helmholtz: results



• $\mu = \rho = 1$, shift $\alpha = 0.2$. Second-order discretization.

- Scaling: largest λ corresponds to Poisson's ratio $\sigma = 0.47$
- Grid size 512×256, 10 grid points per wavelength
- GMRES(5) + 3-level W-cycle as a preconditioner
- Galerkin coarsening

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Multigrid for elastic Helmholtz: results

Shifted Laplacian: acoustic (shear) vs. elastic								
Grid size	400×128		800 >	$\times 256$	1600×512			
ω	2.4π	3.5π	4.7π	7.1π	9.4π	14.2π		
Acoustic	25	40	45	86	75	196		
Elastic	27	37	47	78	79	148		

- Linear heterogeneous media
- Acoustic: Jacobi W(2,2) cycles
- Elastic: Vanka W(1,1) cycles
- Added shift of 0.2 for both
- GMRES(5) + 3-level W-cycle as a preconditioner
- Galerkin coarsening

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Matrix-free geometric MG

The standard 5 point stencil

$$-\mu \Delta_{h} - \omega^{2} M = \frac{\mu}{h^{2}} \begin{bmatrix} -1 & -1 \\ -1 & 4 - \frac{h^{2}}{\mu} \rho \omega^{2} (1 - i\gamma) & -1 \\ -1 & -1 \end{bmatrix}$$

is not efficient enough for re-discretization.

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is not efficient enough for re-discretization. Acoustic: [Singer & Turkel 1998], [Umetani et. al. 2009]

$$\frac{1}{h^2} \begin{bmatrix} -1/6 & -2/3 & -1/6 \\ -2/3 & 10/3 & -2/3 \\ -1/6 & -2/3 & -1/6 \end{bmatrix} - \begin{bmatrix} & -1/12 \\ -1/12 & -2/3 & -1/12 \\ & -1/12 \end{bmatrix} \kappa^2 \omega^2 (1-\gamma \iota)$$

We adapt it to the elastic mixed formulation.

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 β -spread mass matrix:

$$\mathcal{M}^eta =
ho \omega^2 (1-\gamma \imath) \left(eta \left[1
ight] + (1-eta) \cdot rac{1}{4} egin{bmatrix} 1 & 1 \ 1 & 1 \end{bmatrix}
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 β -spread first derivative:

$$\left(\partial_{\mathbf{x}_1}
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 β -spread first derivative:

$$\left(\partial_{x_1}\right)_{h/2}^{\beta} = \frac{1}{h} \left(\beta \begin{bmatrix} -1 & * & 1 \end{bmatrix} + (1-\beta) \cdot \frac{1}{4} \begin{bmatrix} -1 & & 1 \\ -2 & * & 2 \\ -1 & & 1 \end{bmatrix} \right),$$

Divergence and gradient:

$$(\nabla \cdot)_{h}^{\beta} = \begin{pmatrix} (\partial_{x_{1}})_{h/2}^{\beta} & (\partial_{x_{2}})_{h/2}^{\beta} \end{pmatrix}, \quad \nabla_{h}^{\beta} = \begin{pmatrix} (\partial_{x_{1}})_{h/2}^{\beta} \\ (\partial_{x_{2}})_{h/2}^{\beta} \end{pmatrix}$$

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Singer & Turkel's acoustic stencil ($\beta = 2/3$):

$$\nabla_h^T \nabla_h^\beta + M^\beta = (\nabla \cdot)_h^\beta (\nabla \cdot)_h^T + M^\beta = -\Delta_h^\beta + M^\beta.$$



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Discretization for the elastic equation in mixed formulation:

$$\begin{pmatrix} \vec{\nabla}_{h}^{T} A_{e}(\mu) \vec{\nabla}_{h}^{\beta} - M^{\beta} A_{f}(\rho) & (\nabla \cdot)_{h}^{T} \\ (\nabla \cdot)_{h}^{\beta} & \text{diag} \left(\frac{1}{\lambda + \mu} \right) \end{pmatrix} \begin{pmatrix} \vec{\mathbf{u}} \\ \mathbf{p} \end{pmatrix} = \begin{pmatrix} \vec{\mathbf{q}}_{s} \\ \mathbf{0} \end{pmatrix}$$

where A_e is edge averaging, A_f is face averaging.

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Matrix-free geometric MG: results

Tuning the stencil (with $G_s = 10$ and $\sigma = 0.499$):



Damping w = 0.7 near optimal, $\beta = 2/3$ optimal



Minimal shift 0.03 for the spread discretization, 0.11 for standard

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Matrix-free geometric MG: results

Numerical dispersion:



Multigrid convergence on the shifted problem:

LFA two-grid factor and convergence factor in practice									
discretization	$\left \begin{array}{c}\omega = \\\rho_{loc}\end{array}\right $	$\frac{\pi}{5h}, \alpha = c_f$	$= 0.15 \\ \mu_{loc}^2$	$\begin{vmatrix} \omega = \cdot \\ \rho_{loc} \end{vmatrix}$	$\frac{\pi}{4h}, \alpha \in C_f$	$= 0.2 \\ \mu_{loc}^2$	$\begin{array}{l} \omega = \\ \rho_{loc} \end{array}$	$\frac{\pi}{3.3h}, \alpha$ c_f	$\mu = 0.3$ μ_{loc}^2
$\begin{array}{c} \beta = 1 \\ \beta = 2/3 \end{array}$	$\begin{vmatrix} 0.73 \\ 0.37 \end{vmatrix}$	$\begin{array}{c} 0.6 \\ 0.24 \end{array}$	$\begin{array}{c} 0.35 \\ 0.31 \end{array}$	$\begin{vmatrix} 0.81 \\ 0.4 \end{vmatrix}$	$\begin{array}{c} 0.74 \\ 0.27 \end{array}$	$\begin{array}{c} 0.38\\ 0.34 \end{array}$	$\begin{array}{c} 0.75 \\ 0.47 \end{array}$	$\begin{array}{c} 0.7 \\ 0.35 \end{array}$	$\begin{array}{c} 0.44 \\ 0.39 \end{array}$

• The convergence factor is defined by

$$c_f^{(k)} = \left(\frac{\|r_k\|}{\|r_0\|}\right)^{1/k}$$

• μ^2_{loc} serves as a best-case lower bound.

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Matrix-free geometric MG: results



Results for Marmousi geophysical model:

Iteration count for Marmousi-2 elastic media							
	$\beta =$	$= 1, G_s =$	= 10	$\beta = 2/3, G_s = 10(8)$			
	2-level	3-level	4-level	2-level	3-level	4-level	
Grid size	$\alpha = 0.1$	$\alpha = 0.4$	$\alpha = 0.5$	$\alpha = 0.1$	$\alpha=0.3(0.4)$	$\alpha=0.4(0.5)$	
544×112	31	107	147	28(32)	67(106)	100(194)	
1088×224	48	172	238	39(51)	$102 \ (174)$	157 (338)	
2176×448	70	243	374	53(63)	137 (242)	238(572)	

- α is the added attenuation
- G_s grid points per shear wavelength

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Conclusion

- Mixed formulation enables elastic shifted Laplacian MG.
- Our shifted Laplacian MG scales w.r.t. Poisson's ratio.
- Our spread discretization enables a matrix-free method.
- Less shift is needed and less grid points per wavelength.

Future: Better 3D, Block-preconditioning, better acoustic solvers



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- A hybrid shifted Laplacian multigrid and domain decomposition preconditioner for the elastic Helmholtz equations, JCP, [Triester & Y., 2024]

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