# LFA-tuned matrix-free multigrid for the elastic Helmholtz equations

### Racheli Yovel, Eran Treister

Computer Science Department, Ben Gurion University of the Negev, Be'er Sheva, Israel.



The acoustic Helmholtz equation in heterogeneous media

$$
Hp = \rho \nabla \cdot \rho^{-1} \nabla p + \omega^2 \kappa^2 (1 - \gamma \eta) p = q
$$

The acoustic Helmholtz equation in heterogeneous media

$$
Hp = \rho \nabla \cdot \rho^{-1} \nabla p + \omega^2 \kappa^2 (1 - \gamma \eta) p = q
$$

The elastic Helmholtz equation in isotropic heterogeneous media  $\mathcal{H} \vec{u} = \nabla (\lambda + \mu) \nabla \cdot \vec{u} + (\vec{\nabla \cdot} \mu \vec{\nabla}) \vec{u} + \omega^2 \rho (1 - \gamma \eta) \vec{u} = \vec{q}$ Pressure and shear wave velocities:  $V_\rho = \sqrt{\frac{\lambda+2\mu}{\rho}}$  and  $V_s = \sqrt{\frac{\mu}{\rho}}$ .

The acoustic Helmholtz equation in heterogeneous media

$$
Hp = \rho \nabla \cdot \rho^{-1} \nabla p + \omega^2 \kappa^2 (1 - \gamma \eta) p = q
$$

The elastic Helmholtz equation in isotropic heterogeneous media

$$
\mathcal{H}\vec{u} = \nabla(\lambda + \mu)\nabla \cdot \vec{u} + (\vec{\nabla} \cdot \mu \vec{\nabla})\vec{u} + \omega^2 \rho (1 - \gamma \mu \vec{u}) = \vec{q}
$$

Pressure and shear wave velocities:  $V_\rho = \sqrt{\frac{\lambda+2\mu}{\rho}}$  and  $V_s = \sqrt{\frac{\mu}{\rho}}$ .

#### Boundary Conditions

• Neumann / Dirichlet: cause reflections.

The acoustic Helmholtz equation in heterogeneous media

$$
Hp = \rho \nabla \cdot \rho^{-1} \nabla p + \omega^2 \kappa^2 (1 - \gamma \eta) p = q
$$

The elastic Helmholtz equation in isotropic heterogeneous media

$$
\mathcal{H}\vec{u} = \nabla(\lambda + \mu)\nabla \cdot \vec{u} + (\vec{\nabla} \cdot \mu \vec{\nabla})\vec{u} + \omega^2 \rho (1 - \gamma \mu \vec{u}) = \vec{q}
$$

Pressure and shear wave velocities:  $V_\rho = \sqrt{\frac{\lambda+2\mu}{\rho}}$  and  $V_s = \sqrt{\frac{\mu}{\rho}}$ .

#### Boundary Conditions

- Neumann / Dirichlet: cause reflections.
- Absorbing alternatives: Sommerfeld, **ABC**, PML.

#### Difficulties:

• Complex, indefinite and ill-conditioned

### Difficulties:

- Complex, indefinite and ill-conditioned
- High frequencies  $\Rightarrow$  fine meshes  $\Rightarrow$  large matrix

#### Difficulties:

- Complex, indefinite and ill-conditioned
- High frequencies  $\Rightarrow$  fine meshes  $\Rightarrow$  large matrix
- Nearly incompressible case: the grad-div dominates when

$$
\sigma=\frac{\lambda}{2(\lambda+\mu)}\rightarrow\frac{1}{2}
$$

### Difficulties:

- Complex, indefinite and ill-conditioned
- High frequencies  $\Rightarrow$  fine meshes  $\Rightarrow$  large matrix
- Nearly incompressible case: the grad-div dominates when

$$
\sigma = \frac{\lambda}{2(\lambda + \mu)} \to \frac{1}{2}
$$

• Numerical dispersion





### FD Discretizations for acoustic Helmholtz

- 4-th order [Harari & Turkel 1994], [Singer & Turkel 1998]
- 6-th order [Turkel et. al. 2013]
- Non-compact [Dastur & Liau 2020]
- Lower dispersion compact [Liu 2015], [Chen et. al. 2012]
- Rotated grids [Alghamiry et. al. 2022]
- Wavelength adaptive [Xu & Gao 2018]

### Solvers for acoustic Helmholtz

- Domain decomposition [Gander & Zhang 2013], [Hu & Li 2016], [Tausa et. al. 2020], [Daia et. al. 2022], [Claeys, MS2], [Dolean, MS2], [Gong, MS2], [Rees, IP6]
- Shifted Laplacian multigrid [Erlangga et. al. 2006], [Umetani et. al 2009], [Chen et. al. 2012], [Dwarka, MS2], [Chen, MS1]
- FFT [Beylkin 2009], [Osnabrugge 2016], [Wang et. al. 2020]

# Finite difference discretization and available solvers

#### FD Discretizations for elastic Helmholtz

- 2<sup>nd</sup> order staggered (rotated grids) [Virieux 1986]
- Compact nodal [Sketl & Pratt 1998], [Gosselin-Cliche & Giroux 2014]
- 4<sup>th</sup> order staggered [Levander 1988]
- $\bullet$  4<sup>th</sup> order staggered with mass spreading [Li et. al. 2016]

# [Gap]

#### Multigrid adaptations for elastic Helmholtz

- Shifted Laplacian "as is" [Airaksinen et. al. 2009]
- Shifted Laplacian using line-smoothers. [Rizzuti & Mulder 2016]

Multigrid: iterative methods for elliptic PDEs [Brandt 1977]

- Smoothers (i.e. w-Jacobi, GS) reduce oscillatory error modes
- Coarse grid correction (CGC) reduce the smooth modes



Multigrid: iterative methods for elliptic PDEs [Brandt 1977]

- Smoothers (i.e. w-Jacobi, GS) reduce oscillatory error modes
- Coarse grid correction (CGC) reduce the smooth modes
- Isn't CGC enough? Why smoothers?



Multigrid: iterative methods for elliptic PDEs [Brandt 1977]

- Smoothers (i.e. w-Jacobi, GS) reduce oscillatory error modes
- Coarse grid correction (CGC) reduce the smooth modes
- Isn't CGC enough? Why smoothers?



Left image from: "Why Multigrid Methods Are So Efficient" [Yavneh, 2006]

Multigrid: iterative methods for elliptic PDEs [Brandt 1977]



"Why Multigrid Methods Are So Efficient" [Yavneh, 2006]

Two ways to build the coarse operator:

Two ways to build the coarse operator:

• Galerkin coarsening:  $A_H = RA_hP$ .

Two ways to build the coarse operator:

- Galerkin coarsening:  $A_H = RA_hP$ .
- Re-discretization: using the same stencil on coarse grid.

Two ways to build the coarse operator:

- Galerkin coarsening:  $A_H = RA_hP$ .
- Re-discretization: using the same stencil on coarse grid.

Being matrix-free is not (only) about implementation, but about improving complexity [Pazner, IP2].

## Local Fourier analysis

LFA is a predictive tool for analysis of multigrid cycles [Brandt 1977].

• Two grid operator:

$$
TG = S(I - PAH-1RAh)S
$$

• Two grid operator:

$$
TG = S(I - PAH-1RAh)S
$$

• Smoothing factor:

$$
\mu_{\text{loc}} = \max_{\text{high }\theta} (\rho(\tilde{S}(\theta)))
$$

• Two grid operator:

$$
TG = S(I - PAH-1RAh)S
$$

• Smoothing factor:

$$
\mu_{\text{loc}} = \max_{\text{high }\theta} (\rho(\tilde{S}(\theta)))
$$

• Two-grid factor:

$$
\rho_{\text{loc}} = \max_{\text{low } \theta} (\rho(\widetilde{\mathcal{T}G}(\theta))
$$

• Two grid operator:

$$
TG = S(I - PAH-1RAh)S
$$

• Smoothing factor:

$$
\mu_{\textit{loc}} = \max_{\textit{high } \theta} (\rho(\tilde{S}(\theta)))
$$

• Two-grid factor:

$$
\rho_{\text{loc}} = \max_{\text{low } \theta} (\rho(\widetilde{\mathcal{T}G}(\theta)))
$$





# LFA for multiplicative overlapping smoothers

- The amplification factor  $\approx \frac{|\text{error after}|}{|\text{error before}|}$ |error before|
- We need "middle". "before" and "after" are not enough!

[Sivaloganathan 1991], [Maclachlan & Oosterlee 2011], [Rodrigo et. al. 2016], [Treister & Y. 2024]



#### Standard smoothers: unstable for indefinite problems

Standard smoothers: unstable for indefinite problems Coarse grid correction: might amplify error [Elman et. al. 2001]

Standard smoothers: unstable for indefinite problems Coarse grid correction: might amplify error [Elman et. al. 2001]

Shifted Laplacian [Erlangga et. al. 2006]

Use an attenuated matrix  $H_s = H + i\alpha\omega^2 M$  as a preconditioner.

The attenuated system  $H_s\mathbf{x} = \mathbf{r}$  is solvable by multigrid.

**Standard smoothers:** unstable for indefinite problems Coarse grid correction: might amplify error [Elman et. al. 2001]

Shifted Laplacian [Erlangga et. al. 2006]

Use an attenuated matrix  $H_s = H + i\alpha\omega^2 M$  as a preconditioner.

The attenuated system  $H_s\mathbf{x} = \mathbf{r}$  is solvable by multigrid.





Problem: shifted Laplacian is not efficient for elastic Helmholtz.

Problem: shifted Laplacian is not efficient for elastic Helmholtz.

Our adaptations: a new variable  $p = -(\lambda + \mu)\nabla \cdot \vec{u}$ .



• MAC staggered grid finite differences discretization: [McKee et. al. 2008]



• MAC staggered grid finite differences discretization: [McKee et. al. 2008]



• Cell-wise Vanka smoother: invert a  $5 \times 5$  matrix at every cell. [Vanka 1986]



# Multigrid for elastic Helmholtz: results



•  $\mu = \rho = 1$ , shift  $\alpha = 0.2$ . Second-order discretization.

- Scaling: largest  $\lambda$  corresponds to Poisson's ratio  $\sigma = 0.47$
- Grid size  $512\times256$ , 10 grid points per wavelength
- GMRES(5)  $+$  3-level W-cycle as a preconditioner
- Galerkin coarsening



- Linear heterogeneous media
- Acoustic: Jacobi W(2,2) cycles
- Elastic: Vanka W(1,1) cycles
- Added shift of 0.2 for both
- GMRES(5)  $+$  3-level W-cycle as a preconditioner
- Galerkin coarsening

### Matrix-free geometric MG

The standard 5 point stencil

$$
-\mu\Delta_h - \omega^2 M = \frac{\mu}{h^2} \left[ -1 \quad 4 - \frac{h^2}{\mu} \rho \omega^2 (1 - r) \quad -1 \right]
$$

is not efficient enough for re-discretization.

# Matrix-free geometric MG

The standard 5 point stencil

$$
-\mu\Delta_h - \omega^2 M = \frac{\mu}{h^2} \left[ -1 \quad 4 - \frac{h^2}{\mu} \rho \omega^2 (1 - r\gamma) \quad -1 \right]
$$

is not efficient enough for re-discretization. Acoustic: [Singer & Turkel 1998], [Umetani et. al. 2009]

$$
\frac{1}{h^2} \begin{bmatrix} -1/6 & -2/3 & -1/6 \\ -2/3 & 10/3 & -2/3 \\ -1/6 & -2/3 & -1/6 \end{bmatrix} - \begin{bmatrix} -1/12 & -2/3 & -1/12 \\ -1/12 & -2/3 & -1/12 \\ -1/12 & -1/12 \end{bmatrix} \kappa^2 \omega^2 (1 - \gamma I)
$$

We adapt it to the elastic mixed formulation.

 $\beta$ -spread mass matrix:

$$
\mathsf{M}^\beta = \rho \omega^2 (1-\gamma \eta) \left(\beta \left[1\right] + (1-\beta) \cdot \frac{1}{4} \begin{bmatrix} 1 & 1 \\ 1 & 1 \end{bmatrix}\right).
$$

 $\beta$ -spread mass matrix:

$$
\mathsf{M}^\beta = \rho \omega^2 (1-\gamma \prime) \left(\beta \left[1\right] + (1-\beta) \cdot \frac{1}{4} \begin{bmatrix} & 1 & \\ 1 & & 1 \\ & 1 & \end{bmatrix}\right).
$$

 $\beta$ -spread first derivative:

$$
(\partial_{x_1})^{\beta}_{\hbar/2}=\frac{1}{\hbar}\left(\beta\begin{bmatrix}-1 & * & 1\end{bmatrix}+(1-\beta)\cdot\frac{1}{4}\begin{bmatrix}-1 & & 1\\-2 & * & 2\\-1 & & 1\end{bmatrix}\right),
$$

 $\beta$ -spread mass matrix:

$$
\mathsf{M}^\beta = \rho \omega^2 (1-\gamma \eta) \left(\beta \left[1\right] + \left(1-\beta\right) \cdot \frac{1}{4} \begin{bmatrix} & 1 & \\ 1 & & 1 \\ & 1 & \end{bmatrix}\right).
$$

 $\beta$ -spread first derivative:

$$
\left(\partial_{x_1}\right)^\beta_{h/2}=\frac{1}{h}\left(\beta\begin{bmatrix}-1 & * & 1\end{bmatrix}+(1-\beta)\cdot\frac{1}{4}\begin{bmatrix}-1 & & 1\\-2 & * & 2\\-1 & & 1\end{bmatrix}\right),
$$

Divergence and gradient:

$$
\left(\nabla\cdot\right)_{h}^{\beta}=\left(\left(\partial_{x_{1}}\right)_{h/2}^{\beta}\left(\partial_{x_{2}}\right)_{h/2}^{\beta}\right),\quad\nabla_{h}^{\beta}=\left(\begin{matrix}\left(\partial_{x_{1}}\right)_{h/2}^{\beta}\\ \left(\partial_{x_{2}}\right)_{h/2}^{\beta}\end{matrix}\right)
$$





Singer & Turkel's acoustic stencil  $(\beta = 2/3)$ :

$$
\nabla_h^T \nabla_h^{\beta} + M^{\beta} = (\nabla \cdot)_h^{\beta} (\nabla \cdot)_h^T + M^{\beta} = -\Delta_h^{\beta} + M^{\beta}.
$$



Singer & Turkel's acoustic stencil  $(\beta = 2/3)$ :

$$
\nabla_h^T \nabla_h^{\beta} + M^{\beta} = (\nabla \cdot)_h^{\beta} (\nabla \cdot)_h^T + M^{\beta} = -\Delta_h^{\beta} + M^{\beta}.
$$

Discretization for the elastic equation in mixed formulation:

$$
\begin{pmatrix}\n\vec{\nabla}_{h}^{T} A_{e}(\mu) \vec{\nabla}_{h}^{\beta} - M^{\beta} A_{f}(\rho) & (\nabla \cdot)_{h}^{T} \\
(\nabla \cdot)_{h}^{\beta} & \text{diag}\left(\frac{1}{\lambda + \mu}\right)\n\end{pmatrix}\n\begin{pmatrix}\n\vec{\mathbf{u}} \\
\mathbf{p}\n\end{pmatrix} = \begin{pmatrix}\n\vec{\mathbf{q}}_{s} \\
\mathbf{0}\n\end{pmatrix}
$$

where  $A_e$  is edge averaging,  $A_f$  is face averaging.



Singer & Turkel's acoustic stencil  $(\beta = 2/3)$ :

$$
\nabla_h^T \nabla_h^{\beta} + M^{\beta} = (\nabla \cdot)_h^{\beta} (\nabla \cdot)_h^T + M^{\beta} = -\Delta_h^{\beta} + M^{\beta}.
$$

Discretization for the elastic equation in mixed formulation:

$$
\begin{pmatrix}\n\vec{\nabla}_{h}^{T} A_{e}(\mu) \vec{\nabla}_{h}^{\beta} - M^{\beta} A_{f}(\rho) & (\nabla \cdot)_{h}^{T} \\
(\nabla \cdot)_{h}^{\beta} & \text{diag}\left(\frac{1}{\lambda + \mu}\right)\n\end{pmatrix}\n\begin{pmatrix}\n\vec{\mathbf{u}} \\
\mathbf{p}\n\end{pmatrix} = \begin{pmatrix}\n\vec{\mathbf{q}}_{s} \\
\mathbf{0}\n\end{pmatrix}
$$

where  $A_e$  is edge averaging,  $A_f$  is face averaging.

# Matrix-free geometric MG: results

Tuning the stencil (with  $G_s = 10$  and  $\sigma = 0.499$ ):



Damping  $w = 0.7$  near optimal,  $\beta = 2/3$  optimal



Minimal shift 0.03 for the spread discretization, 0.11 for standard

### Matrix-free geometric MG: results

#### Numerical dispersion:



Multigrid convergence on the shifted problem:



• The convergence factor is defined by

$$
c_f^{(k)} = \left(\frac{\|r_k\|}{\|r_0\|}\right)^{1/k}
$$

•  $\mu_{loc}^2$  serves as a best-case lower bound.

# Matrix-free geometric MG: results



#### Results for Marmousi geophysical model:



- $\alpha$  is the added attenuation
- $G_s$  grid points per shear wavelength

# Conclusion

- Mixed formulation enables elastic shifted Laplacian MG.
- Our shifted Laplacian MG scales w.r.t. Poisson's ratio.
- Our spread discretization enables a matrix-free method.
- Less shift is needed and less grid points per wavelength.

Future: Better 3D, Block-preconditioning, better acoustic solvers



# Thanks

#### Funding:

- The U.S. Department of Energy (travel grant)
- Israel Science foundation (grant No. 1589/19)
- Ariane de Rothschild woman doctoral program
- Kreitman High-tech scholarship, BGU

#### References:

- LFA-tuned matrix-free multigrid method for the elastic Helmholtz equation, SISC, [Y. & Treister, 2024]
- A hybrid shifted Laplacian multigrid and domain decomposition preconditioner for the elastic Helmholtz equations, JCP, [Triester & Y., 2024]

yovelr@bgu.ac.il