# Preconditioning finite difference matrices from density functional theory

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2 A method of efficient preconditioning



#### A family of block linear systems

A method of efficient preconditioning

3 A pair of welcome surprises

# Computing the functional trace of a linear operator

• We want to approximate

$$\int_0^\infty \operatorname{Tr}[f(\mathcal{A}(\omega))] \mathrm{d}\omega$$

by computing the trace at quadrature points  $\omega_k$ 

- Any method we use to approximate the trace will require performing *A*(ω<sub>k</sub>)*V* for a block of vectors *V* (e.g., subspace iteration, Arnoldi quadrature, Hutchinson trace estimator, ...)
- $\mathcal{A}(\omega_k)$  cannot be constructed explicitly or applied in some matrix-free fashion; instead, applying  $\mathcal{A}(\omega_k)$  requires an expensive and complicated sequence of operations

The multiplication  $\mathcal{A}(\omega_k)V$  is performed in three stages:

$$(H - \lambda_j I - i\omega_k I) Y_j = -V \odot \Psi_j, \quad j = 1, 2, ..., n_s$$
  
 $W = 4 \operatorname{Re} \left( \sum_{j=1}^{n_s} \Psi_j^* \odot Y_j \right)$   
 $\nabla^2 (\mathcal{A}(\omega_k) V) = -4\pi W$ 

Stage 1 is the most computationally intensive and requires solving  $n_s$  block linear systems.

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## Hamiltonian matrix (from real-space DFT)

$$\underbrace{(H-\lambda_j I-i\omega_k I)}_A Y_j = -V \odot \Psi_j, \quad j=1,2,\ldots,n_s$$

• 
$$H = -\frac{1}{2}\nabla^2 + \chi\chi^H + \text{diag}(P_{\text{eff}})$$

- *H* is real symmetric and indefinite; its lowest n<sub>s</sub> eigenpairs (λ<sub>j</sub>, Ψ<sub>j</sub>) are known a priori
- A is not Hermitian but satisfies  $A = A^T$  (i.e., it is complex symmetric)
- A becomes nearly singular as  $\omega_k \rightarrow 0$
- We use the convention

$$\lambda_1 < \lambda_2 < \ldots < \lambda_{n_s}$$
$$0 < \omega_8 < \omega_2 < \ldots < \omega_1$$

Algorithm 1 Block preconditioned  $COCG^{a}$  for solving AX = B satisfying  $A = A^T$  and using a preconditioner  $M^{-1}$ 1: Initialize  $X_0$  and  $V_0 \leftarrow B - AX_0$ 2:  $W_0 \leftarrow M^{-1}V_0$ 3:  $\rho_0 \leftarrow V_0^T W_0$ 4:  $P_{-1} \leftarrow 0, \beta_{-1} \leftarrow 0$ 5: for  $j = 0, 1, 2, \dots$  do 6:  $P_i \leftarrow V_i + P_{i-1}\beta_{i-1}$ 7:  $U_j \leftarrow AP_j$ 8:  $\mu_i \leftarrow U_i^T P_i$  $\alpha_i \leftarrow \mu_i^{-1} \rho_i$ 9:  $X_{i+1} \leftarrow X_i + P\alpha_i$ 10: 11:  $V_{i+1} \leftarrow V_i - U_i \alpha_i$ 12:  $W_{i+1} \leftarrow M^{-1}V_{i+1}$ 13:  $\rho_{i+1} \leftarrow V_{i+1}^T W_{i+1}$  $\beta_j \leftarrow \rho_i^{-1} \rho_{j+1}$ 14: 15: end for

- COCG is not optimal in the residual norm but achieves similar convergence to GMRES in practice
- There is potential for breakdown in the iteration

<sup>a</sup>van der Vorst and Melissen 1990.

### The linear systems vary widely in difficulty



## Should we use a block size greater than 1?



We know the lowest  $n_s$  eigenpairs  $(\lambda_\ell, \Psi_\ell)$  of H; since H and A share eigenvectors, we can deflate an invariant subspace from the initial residual:

$$X_{0} = \boxed{\Psi(\Lambda - \lambda_{j}I - i\omega_{k}I)^{-1}\Psi^{H}B}$$

$$R_{0} = B - AX_{0}$$

$$= B - (H - \lambda_{j}I - i\omega_{k}I)\Psi(\Lambda - \lambda_{j}I - i\omega_{k}I)^{-1}\Psi^{H}B$$

$$= B - \Psi\Psi^{H}B$$

$$\Psi^{H}R_{0} = \Psi^{H}B - \Psi^{H}\Psi\Psi^{H}B$$

$$= 0$$

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### Iterations required is also RHS-dependent



We want a preconditioner that reduces the total time required to solve all of the Sternheimer equations. This preconditioner must either

be extremely cheap to apply

or

drastically accelerate convergence

or

ideally both!

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Assume  $M = M^T$ . It admits an *LDU* decomposition:

$$M = LDU$$
 and  $M^T = U^T D^T L^T = U^T DL$ 

By induction, we can show  $U = L^T$  and therefore:

$$M = LDL^{T} = LD^{1/2}D^{1/2}L^{T} = (LD^{1/2})(LD^{1/2})^{T} = KK^{T}$$

The preconditioned matrix in COCG is  $A_{\text{precond}} = K^{-1}AK^{-T}$ , thus:

$$A_{\text{precond}}^{T} = K^{-1}A^{T}K^{-T} = A_{\text{precond}}$$

# Choosing a good preconditioner

$$\left(-\frac{1}{2}\nabla^2 + \chi \chi^{\mathcal{H}} + \mathsf{diag}\left(\mathcal{P}_{\mathsf{eff}}\right) - \lambda_j \mathcal{I} - i\omega_k \mathcal{I}\right) Y_j = -\mathcal{V} \odot \Psi_j$$

The natural choice of preconditioner is

$$M^{-1} = \left(-\frac{1}{2}\nabla^2 + \overline{P_{\text{eff}}}I - \lambda_j I - i\omega_k I\right)^{-1}$$

If we know the eigendecomposition  $abla^2 = Q \Lambda Q^H$ ,  $M^{-1}$  simply becomes

$$M^{-1} = Q \left( -\frac{1}{2} \Lambda + (\overline{P_{\text{eff}}} - \lambda_j - i\omega_k) I \right)^{-1} Q^H$$

$$\nabla^{2} = (\mathcal{I}_{z} \otimes \mathcal{I}_{y} \otimes \mathcal{D}_{x}) + (\mathcal{I}_{z} \otimes \mathcal{D}_{y} \otimes \mathcal{I}_{x}) + (\mathcal{D}_{z} \otimes \mathcal{I}_{y} \otimes \mathcal{I}_{x})$$
  
If  $\mathcal{D}_{x}$ ,  $\mathcal{D}_{y}$ , and  $\mathcal{D}_{z}$  have eigenvectors  $Q_{x}$ ,  $Q_{y}$ , and  $Q_{z}$ , respectively, then:  
$$Q = Q_{z} \otimes Q_{y} \otimes Q_{x} \text{ and } Q^{H} = Q_{z}^{H} \otimes Q_{y}^{H} \otimes Q_{x}^{H}$$
$$M^{-1}u = Q\Lambda_{M^{-1}}Q^{H}u = (Q_{z} \otimes Q_{y} \otimes Q_{x})\Lambda_{M^{-1}}(Q_{z}^{H} \otimes Q_{y}^{H} \otimes Q_{x}^{H})u$$

The product  $(A \otimes B \otimes C)u$  can be computed efficiently *even when the finite difference directions are non-orthogonal*<sup>1</sup>.

<sup>&</sup>lt;sup>1</sup>Sharma and Suryanarayana 2018.

# Spectral behavior: easy Sternheimer system $(\lambda_1, \omega_1)$



# Spectral behavior: hardest Sternheimer system $(\lambda_{32}, \omega_8)$



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### Block size dependence returns with a twist



## Uniform difficulty across the preconditioned systems



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## When should we use this preconditioner?



## Modest speedup even before optimizing



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- Preconditioning reintroduces the block size dependence and is faster when large block sizes are used on the hardest linear systems
- Shockingly, preconditioning fixes the issue of load imbalance that cropped up due to our specific right-hand side vectors

# Full comparison for $(\lambda_8, \omega_8)$

